



# Third Moment Description of the Turbulent Cascade and Intermittency

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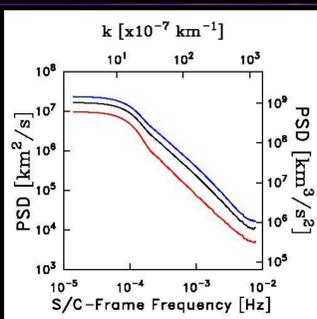
## Magnetohydrodynamic Energy Cascade Rate

The energy dissipation rate of an incompressible magnetohydrodynamic [MHD] fluid was derived by Politano & Pouquet [1] following Yaglom [2] by assuming statistical homogeneity and strict spatial independence. Here, we present the isotropic form after assuming the existence of an inertial range and the second moment to be time stationary:

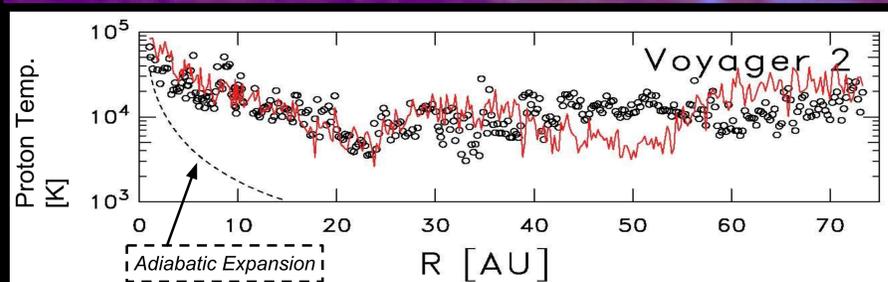
$$\frac{4}{3}\epsilon^\pm r_L = -\langle \delta z_L^\mp(\mathbf{r}) |\delta z_i^\pm(\mathbf{r})|^2 \rangle$$

The scale invariant energy cascade rate is related to a particular third order structure function through the spatial increment.

The Kolmogorov phenomenology predicts the power law and offers a universal description of the small scale structure of turbulence. **Right)** The power spectral density is defined:  $E^{\text{tot}}(\mathbf{k})$ ,  $E^{\text{out}}(\mathbf{k})$ ,  $E^{\text{in}}(\mathbf{k})$ . The outward fluctuations possess more energy, all three possess a exponent of  $-5/3$  through the inertial range.

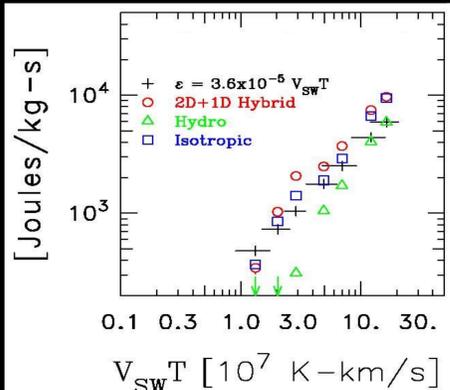


## Proton Heating Rate in the Solar Wind at 1AU



**Figure Above)** The non-adiabatic radial temperature dependence of solar wind protons (circles) requires in situ heating and pickup ions (red line) [cite].

**Figure Right)** For a large data set the third moment structure function converges to a scale independent energy cascade rate. The average of the measured cascade rates match the inferred heating rate from the velocity-temperature product. [3,4,5]



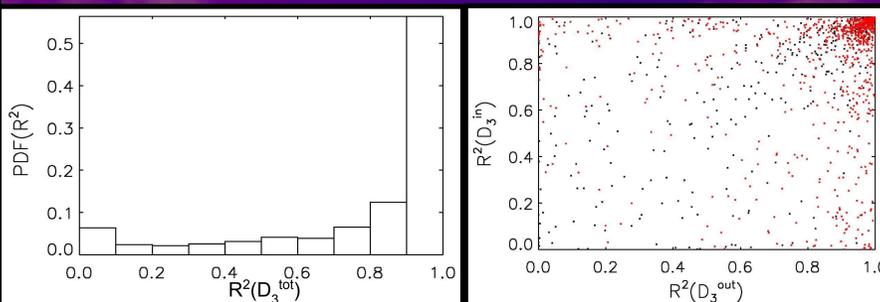
## Variable Cascade Rates and Intermittency

Traditional studies of intermittency account for the non-unique exponents of the structure functions, or measure the kurtosis of the field fluctuations. Our studies measure local dynamics to present a new picture; local inertial range dynamics are well behaved and adjust as large-scale and dissipation scales are not in statistical equilibrium.

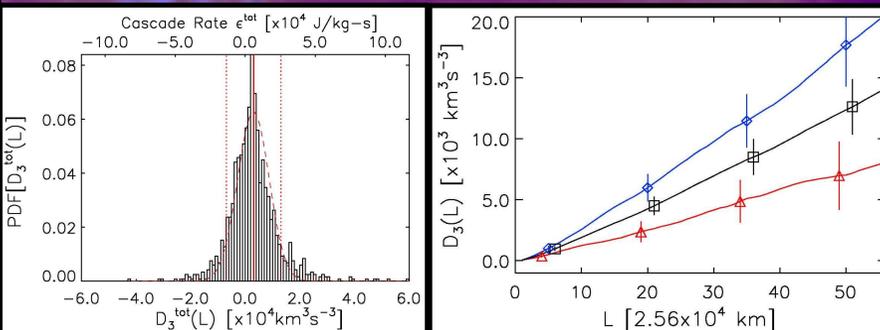
The averaging scheme employed in our studies is two-tiered:

- I. Calculate the third moment function for a few correlation lengths [Local Energy Cascade Rate].
- II. Interpolate the functions to common lags and then calculate Gaussian statistics [Heating Rate].

**Figures Below)** The energy range (1200 - 2800 [km<sup>2</sup>/s<sup>2</sup>]) and normalized cross helicity range (-0.75 - 0.75) was chosen to study the linearity ( $R^2 = \text{least-squares}$ ) of the third moment function. **Left)** More than 70% of the local estimates are linear. **Right)** Red points are  $R^2(D_3^{\text{tot}}) > 0.7$ . Linear  $D_3^{\text{tot}}$  are often accompanied by linear  $D_3^{\text{out}}$  and  $D_3^{\text{in}}$ .

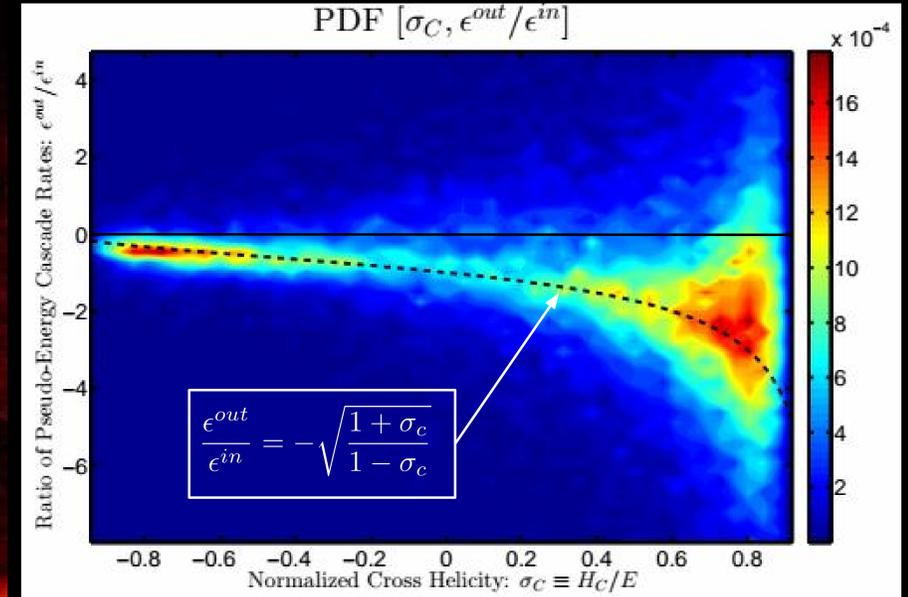


**Figures Below)** Only  $R^2(D_3^{\text{tot}}) > 0.7$  are included. **Left)** Many samples produce a well formed distribution and positive average. **Right)** 3<sup>rd</sup> moment functions estimate a scale independent cascade rate.  $D_3^{\text{tot}}$ ,  $D_3^{\text{out}}$ ,  $D_3^{\text{in}}$ .



The energy cascade rates are being determined on a local time scale and possess a large degree of variability. Energy transfer in the inertial range displays a large degree of back transfer. The heating rate can be measured from the average of many local energy cascade rates [6].

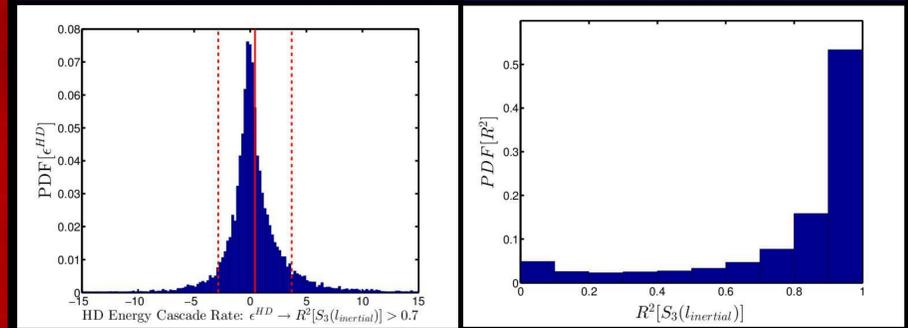
## Local Cascade Rates and Anti-Correlation



**Figure Above)** The ability of the third moment technique to estimate the local energy cascade rate can then be employed to study various decay processes and regions of the solar wind [7].

## Hydrodynamic Cascade and Intermittency

The measurement [Modane Wind Tunnel] of the hydrodynamic energy cascade rate produces very similar results to the solar wind. The inertial range bounds are not in statistical equilibrium on similar time-scales. This creates a very dynamic picture of the energy transfer in the inertial range bound by local quantities [8].



## References and Acknowledgements

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[8] Coburn et al. (In Press)  
Background Figure: SolarDynamicObservatory/NASA